

Ph135 - Problem Set 1

October 21, 2018

Problem 1

(a)

$$U = \int_0^\infty \rho_0 \epsilon \frac{1}{e^{\beta(\epsilon-\mu)} + 1} d\epsilon \quad (1)$$

(b)

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V, \quad (2)$$

Now by the Sommerfeld expansion,

$$U = U_0 + \frac{\pi^2}{6} k_B^2 T^2 \rho_0. \quad (3)$$

This is expected, because the only place that electrons move around is T about the Fermi surface, and when they move around, they are excited by an amount T , so we have $U \propto T^2$. Therefore,

$$C_V = \frac{\pi^2}{3} \rho_0 k_B T \quad (4)$$

Problem 2

(a)

$$\rho(E) dE = \rho(\epsilon_k) d\epsilon_k, \quad (5)$$

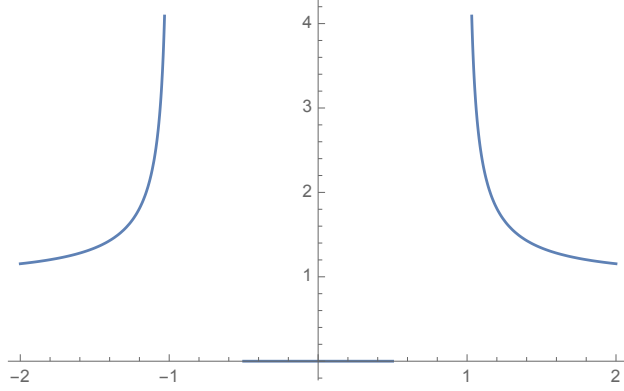
which implies that

$$\rho(E) = \rho_0 \left(\frac{dE}{d\epsilon_k} \right)^{-1}. \quad (6)$$

We plug in the expression in (1), and get that

$$\rho(E) = \rho_0 \frac{|E|}{\sqrt{E^2 - \Delta^2}} \quad (7)$$

when $|E| \geq \Delta$. $\rho(E) = 0$ when $|E| < \Delta$, because $|E|$ cannot be smaller than Δ .



(b) The plot of density of states looks similar to Fig. 1.

(c)

$$U = \int dE \rho(E) \frac{E}{e^{\beta(E-\mu)} + 1} = \int_{-\infty}^{-\Delta} dE \rho_0(E) \frac{-E^2}{\sqrt{E^2 - \Delta^2}} \frac{1}{e^{\beta(E-\mu)} + 1} \quad (8)$$

$$+ \int_{\Delta}^{\infty} dE \rho_0(E) \frac{E^2}{\sqrt{E^2 - \Delta^2}} \frac{1}{e^{\beta(E-\mu)} + 1} \quad (9)$$

(d) We do a change of variable $x \equiv \sqrt{E^2 - \Delta^2}$, then

$$U = \int_0^{\infty} dx \rho_0(x) \sqrt{x^2 + \Delta^2} \frac{1}{e^{\beta(\sqrt{x^2 + \Delta^2})} + 1}. \quad (10)$$

When T is much less than Δ , we have $\epsilon_{\vec{k}} \sim E_F$, we can expand U for small x , we get

$$\begin{aligned} U &\approx \int_0^{\epsilon} dx \rho_0(x) \left(\Delta + \frac{x^2}{2\Delta} \right) e^{-\beta(\Delta + x^2/(2\Delta))} \\ &\sim \Delta^{3/2} T^{1/2} e^{-\Delta/T}, \end{aligned} \quad (11)$$

and

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V \sim \Delta^{5/2} T^{-3/2} e^{-\Delta/T}. \quad (12)$$

We can see that because the lowest energy excitation has energy Δ the heat capacity is suppressed exponentially with the Boltzmann factor.

Problem 3

At $T = 0$, the total energy is given by

$$\begin{aligned} U &= 2 \times \frac{1}{8} \times \int_0^{k_F} 4\pi k^2 \frac{V dk}{(2\pi)^3} \frac{\hbar^2 k^2}{2m} \\ &= \frac{\hbar^2 k_F^5 V}{80\pi^2 m}. \end{aligned} \quad (13)$$

Now

$$N = 2 \times \frac{1}{8} \times \frac{4}{3} \times \frac{k_F^3 V}{(2\pi)^3} = \frac{k_F^3 V}{24\pi^3}, \quad (14)$$

Therefore,

$$U = \frac{3}{5} N \frac{\hbar^2 k_F^2}{2m} = \frac{3}{5} N \frac{\hbar^2}{2m} \left(\frac{3\pi^2 N}{V} \right)^{2/3}, \quad (15)$$

and

$$P = -\frac{\partial U}{\partial V} = \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \left(\frac{N}{V} \right)^{5/3} = \frac{(3\pi^2)^{2/3} \hbar^2}{5m} n^{5/3}. \quad (16)$$

For example, for copper, $n \approx 8.5 \times 10^{28}/m^3$, therefore,

$$P \approx 1.91 \times (1.05 \times 10^{-34})^2 / (9 \times 10^{-31}) \times (8.5 \times 10^{28})^{5/3} \text{ Pa} \approx 3.85 \times 10^{10} \text{ Pa}. \quad (17)$$